

Explosive solutions of semilinear elliptic systems with gradient term

Marius Ghergu and Vicențiu Rădulescu

Abstract. We study the existence of boundary blow-up solutions to the nonlinear elliptic system $\Delta u + |\nabla u| = p(|x|)f(v), \Delta v + |\nabla v| = q(|x|)g(u)$ in Ω . Here Ω is either a bounded domain in \mathbb{R}^N or it denotes the whole space. The nonlinearities f and g are positive and continuous, while the nonnegative potentials p and q are continuous and satisfy appropriate growth conditions at infinity. We show that boundary blow-up positive solutions fail to exist if f and g are sublinear. This result holds both if Ω is bounded, and if Ω is the whole space but p and q have slow decay at infinity. We establish the existence of infinitely many entire blow-up solutions in the case where p and q are of fast decay and if f and g satisfy a sublinear type growth condition at infinity.

Soluciones explosivas de sistemas elípticos semilineales con términos gradientes

Resumen. Estudiamos la existencia de soluciones del sistema elíptico no lineal $\Delta u + |\nabla u| = p(|x|)f(v), \Delta v + |\nabla v| = q(|x|)g(u)$ en Ω que explotan en el borde. Aquí Ω es un dominio acotado de \mathbb{R}^N o el espacio total. Las nolinealidades f y g son funciones continuas positivas mientras que los potenciales p y q son funciones continuas que satisfacen apropiadas condiciones de crecimiento en el infinito. Demostramos que las soluciones explosivas en el borde dejan de existir si f y g son sublineales. Esto se tiene o bien si Ω es acotado o cuando Ω es el espacio total pero p y q decaen lentamente en el infinito. Mostramos la existencia de infinitas soluciones enteras explosivas cuando p y q decaen rápidamente y cuando f y g satisfacen una condición de tipo sublineal en el infinito.

1. Introduction and the main results

Existence and nonexistence of solutions of the semilinear elliptic system

$$\begin{cases} \Delta u = f(x, u, v) & \text{in } \Omega, \\ \Delta v = g(x, u, v) & \text{in } \Omega \end{cases}$$
(1)

have received much attention recently. See, for example, Chen and Lu [2], Cîrstea and Rădulescu [4], Clément, Manásevich and Mitidieri [5], Dalmasso [6], De Figueiredo and Jianfu [7], Lair and Shaker [14], Serrin and Zou [18, 19], Yarur [20], Wang and Wood [21], and the references therein. Most of these results have to do with the nonexistence of positive solutions, the existence of radial solutions, or the asymptotic behavior of solutions.

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We are concerned in this paper with the study of positive solutions to the following class of semilinear elliptic systems with gradient term

$$\begin{cases} \Delta u + |\nabla u| = p(|x|)f(v) & \text{in }\Omega, \\ \Delta v + |\nabla v| = q(|x|)g(u) & \text{in }\Omega, \end{cases}$$
(2)

where $\Omega \subset \mathbb{R}^N$ $(N \ge 1)$ denotes either a bounded open set in \mathbb{R}^N or the whole of \mathbb{R}^N . Throughout this paper we assume that $p, q \neq 0$ are nonnegative Hölder functions. We also assume that f and g are Hölder, positive and non-decreasing functions on $(0, \infty)$.

We are mainly interested in finding properties of *large (explosive, blow-up) solutions* of (2), that is positive solutions (u, v) satisfying $u(x) \to +\infty$ and $v(x) \to +\infty$ as dist $(x, \partial\Omega) \to 0$ (if Ω is bounded), or $u(x) \to +\infty$ and $v(x) \to +\infty$ as $|x| \to \infty$ (if $\Omega = \mathbb{R}^N$). In the latter case such solutions are called *entire large (explosive, blow-up) solutions*. A geometric motivation in that sense can be found in [3, 12, 15]. We also point out the pioneering work of Keller [10] and Osserman [16].

The corresponding equation that leads us to the system (2) is

$$\Delta u + |\nabla u|^a = p(x)f(u), \quad x \in \Omega, \ 0 < a \le 2,$$

which was treated in [1, 8] (in the case where Ω is bounded) and in [9, 13] (for $\Omega = \mathbb{R}^N$). Problems of this type arise in stochastic control theory and have been first studied in Lasry and Lions [11]. The corresponding parabolic equation was considered in Quittner [17]. In terms of the dynamic programming approach, an explosive solution of (2) corresponds to a value function (or Bellman function) associated to an infinite exit cost (see [11]).

Our first result asserts that if Ω is bounded and if both f and g are sublinear at infinity, then problem (2) has no positive boundary blow-up solution. More precisely, the following hold

Theorem 1 Suppose $\Omega \subset \mathbb{R}^N$ is a bounded domain and f, g satisfy

$$\max\left\{\sup_{t\geq 1}\frac{f(t)}{t}, \ \sup_{t\geq 1}\frac{g(t)}{t}\right\} < +\infty.$$
 (A₁)

Then problem (2) has no positive large solution.

The same conclusion holds if $\Omega = \mathbb{R}^N$, but under natural additional assumptions related to the behavior of p and q at infinity. In order to state the result in this case, let us first define, for any $r \ge 0$,

$$P(r) = \frac{\int_0^r e^t t^{N-1} p(t) dt}{e^r r^{N-1}}, \quad Q(r) = \frac{\int_0^r e^t t^{N-1} q(t) dt}{e^r r^{N-1}}.$$
(3)

Theorem 2 Let $\Omega = \mathbb{R}^N$. Assume that (A_1) holds and

$$\int_{1}^{\infty} P(r)dr < +\infty, \quad \int_{1}^{\infty} Q(r)dr < +\infty.$$
(4)

Then problem (2) has no positive entire large solution.

Theorem 3 Let $\Omega = \mathbb{R}^N$. Assume that

$$\int_{1}^{\infty} P(r)dr = +\infty, \quad \int_{1}^{\infty} Q(r)dr = +\infty.$$
(5)

If

$$\lim_{t \to \infty} \frac{f(ag(t))}{t} = 0, \quad \text{for all constants} \ a \ge 1, \tag{A2}$$

then problem (2) has infinitely many positive entire large solutions.

We point out that Condition (A_2) has been introduced in [4].

Remark 1 Using the fact that

$$\int_0^r e^r t^k dt = k! e^r \sum_{s=1}^k (-1)^{k-s} \frac{t^s}{s!} \quad \text{for all integers } k \ge 1,$$
(6)

we observe that the following functions verify (4) or (5): (i) condition (4) holds provided that $p(t) = \frac{1}{1+t^{\gamma}}$, $\gamma > 1$ and $q(t) = \frac{1}{(1+t^2)^{\theta}}$, $\theta > \frac{1}{2}$. (ii) condition (5) holds provided that $p(t) = t^{\gamma}, \ q(t) = t^{\theta}, \ \gamma, \theta \ge 0.$

Remark 2 We give in what follows some examples of nonlinearities f and g that satisfy (A_2) : (i) $f(t) = \sum_{j=1}^{l} a_j t^{\gamma_j}$, $g(t) = \sum_{k=1}^{m} b_k t^{\theta_k}$, $t \ge 0$ with $a_j, b_k, \gamma_j, \theta_k > 0$ and $\gamma \theta < 1$, where $\gamma = \max_{1 \le j \le l} \gamma_j$, $\theta = \max_{1 \le k \le m} \theta_k$. (ii) $f(t) = (1 + t^{\gamma_1})^{\gamma_2}$, $g(t) = (1 + t^{\theta_1})^{\theta_2}$, where $\gamma_1, \gamma_2, \theta_1, \theta_2 > 0$ and $\gamma_1 \gamma_2 \theta_1 \theta_2 < 1$. (iii) $f(t) = \ln(1 + t^{\gamma})$, $g(t) = \ln(1 + t^{\theta})$, $\gamma, \theta > 0$. (iv) $f(t) = \ln(1+t^{\gamma}), g(t) = e^{t^{\theta}}, \gamma > 0, \theta \in (0,1).$

Proof of Theorem 1 2.

Suppose that (u, v) is a positive large solution of (2) and let $w(x) = \ln(1 + u(x) + v(x)), x \in \Omega$. Then w is a positive function and $w(x) \to \infty$ as dist $(x, \partial \Omega) \to 0$. A simple calculation yields

$$\Delta w = \frac{\Delta u + \Delta v}{1+u+v} - \frac{\displaystyle\sum_{i=1}^{N}(u_{x_i}+v_{x_i})^2}{(1+u+v)^2} \quad \text{in } \Omega$$

Taking into account the assumption (A_1) we have

$$\begin{aligned} \Delta w &\leq \frac{\Delta u + \Delta v}{1 + u + v} \\ &\leq \frac{\|p\|_{L^{\infty}(\Omega)} f(v) + \|q\|_{L^{\infty}(\Omega)} g(u)}{1 + u + v} \\ &\leq \left(\|p\|_{L^{\infty}(\Omega)} + \|q\|_{L^{\infty}(\Omega)}\right) \frac{f(v) + g(u)}{1 + u + v} \\ &\leq \left(\|p\|_{L^{\infty}(\Omega)} + \|q\|_{L^{\infty}(\Omega)}\right) \left(\frac{f(1 + v)}{1 + v} + \frac{g(1 + u)}{1 + u}\right) \leq K, \end{aligned}$$

for some constant K > 0. Hence

 $\Delta(w(x) - K|x|^2) < 0, \quad \text{for all } x \in \Omega.$

Let $z(x) = w(x) - K|x|^2, x \in \Omega$. Then

$$\Delta z < 0 \quad \text{in } \Omega \tag{7}$$

and

$$z(x) \to \infty$$
 as $\operatorname{dist}(x, \partial \Omega) \to 0.$ (8)

Fix $x_0 \in \Omega$ and M > 0. At this point, to reach a contradiction we will show that $z(x_0) > M$. Suppose $z(x_0) \leq M$. For all $\delta > 0$, we set

$$\Omega_{\delta} = \{ x \in \Omega \mid \operatorname{dist}(x, \partial \Omega) > \delta \}$$

Since $z(x) \to \infty$ as $\operatorname{dist}(x, \partial \Omega) \to 0$, we can choose $\delta > 0$ such that z(x) > M for all $x \in \Omega \setminus \Omega_{\delta}$. Obviously, $x_0 \in \Omega_{\delta}$. Moreover, $M - z(x_0) \ge 0$ and $(M - z)|_{\partial \Omega_{\delta}} \le 0$. Therefore we can find $\bar{x} \in \Omega_{\delta}$ such that

$$\max_{\overline{\Omega}_{\delta}}(M - z(x)) = M - z(\bar{x}) \le 0.$$

It follows that $\Delta(M - z)(\bar{x}) \leq 0$, that is $\Delta z(\bar{x}) \geq 0$ which contradicts (7). Hence (2) has no positive large solutions. This completes the proof.

Remark 3 We can employ the same method as above to show that the system

$$\begin{cases} \Delta u + |\nabla v| = p(|x|)f(v) & \text{ in } \Omega, \\ \Delta v + |\nabla u| = q(|x|)g(u) & \text{ in } \Omega, \end{cases}$$

has no positive large solutions if f and g satisfy (A_1) .

3. Proof of Theorem 2

Arguing by contradiction, let us assume that the system (2) has the positive entire large solution (u, v). Consider the spherical average of u and v defined by

$$\bar{u}(r) = \frac{1}{c_N r^{N-1}} \int_{|x|=r} u(x) d\sigma_x, \quad r \ge 0$$
(9)

$$\bar{v}(r) = \frac{1}{c_N r^{N-1}} \int_{|x|=r} v(x) d\sigma_x, \quad r \ge 0$$
(10)

where c_N is the surface area of the unit sphere in \mathbb{R}^N . Since u and v are positive entire large solutions it follows that \bar{u}, \bar{v} are positive and $\lim_{r \to \infty} \bar{u}(r) = \lim_{r \to \infty} \bar{v}(r) = +\infty$. By the change of variable $x \to ry$, we have

$$\bar{u}(r) = \frac{1}{c_N} \int_{|y|=1} u(ry) \, d\sigma_y, \quad r \ge 0$$

and

$$\bar{u}'(r) = \frac{1}{c_N} \int_{|y|=1} \nabla u(ry) \cdot y \, d\sigma_y, \quad r \ge 0.$$
(11)

The above relation may be rewritten as

$$\bar{u}'(r) = \frac{1}{c_N} \int_{|y|=1} \frac{\partial u}{\partial r} (ry) \, d\sigma_y = \frac{1}{c_N r^{N-1}} \int_{|x|=r} \frac{\partial u}{\partial r} (x) \, d\sigma_x,$$

that is

$$\bar{u}'(r) = \frac{1}{c_N r^{N-1}} \int_{|x|=r} \Delta u(x) \, d\sigma_x, \quad \text{for all } r \ge 0.$$
(12)

Similarly we have

$$\bar{v}'(r) = \frac{1}{c_N r^{N-1}} \int_{|x|=r} \Delta v(x) \, d\sigma_x, \quad \text{for all } r \ge 0.$$
(13)

Due to the presence of the gradient term in (2), we cannot infer that $\Delta u \ge 0$ in \mathbb{R}^N and so we do not know if $\bar{u}' \ge 0$ (or $\bar{v}' \ge 0$) in $[0, \infty)$. In order to overcome this lack of monotonicity, set

$$U(r) = \max_{0 \le t \le r} \bar{u}(r), \quad V(r) = \max_{0 \le t \le r} \bar{v}(r).$$
(14)

Now it is easy to see that U, V are positive and non-decreasing functions. Moreover $U \ge \overline{u}, V \ge \overline{v}$ and $U(r), V(r) \to +\infty$ as $r \to \infty$.

By (A_1) , that there exists M > 0 such that

$$\max\{f(t), g(t)\} \le M(1+t), \text{ for all } t \ge 0.$$
(15)

Now (11), (12) and (15) lead to

$$\begin{split} \bar{u}'' + \frac{N-1}{r} \, \bar{u}' + \bar{u}' &\leq \frac{1}{c_N r^{N-1}} \int_{\substack{|x|=r\\ |x|=r}} [\Delta u(x) + |\nabla u|(x)] \, d\sigma_x \\ &= p(r) \frac{1}{c_N r^{N-1}} \int_{\substack{|x|=r\\ |x|=r}} f(v(x)) \, d\sigma_x \\ &\leq M p(r) \frac{1}{c_N r^{N-1}} \int_{\substack{|x|=r\\ |x|=r\\ |x|=r}} (1+v(x)) \, d\sigma_x \\ &= M p(r) \left(1 + \bar{v}(r)\right) \\ &\leq M p(r) \left(1 + V(r)\right), \end{split}$$

for all $r \ge 0$. It follows that

$$(r^{N-1}e^r \bar{u}')' \le M e^r r^{N-1} p(r) (1+V(r)) \text{ for all } r \ge 0.$$

So, for all $r \ge r_0 > 0$,

$$\begin{split} \bar{u}(r) &\leq \bar{u}(r_0) + M \int_{r_0}^r e^{-t} t^{1-N} \int_0^t e^s s^{N-1} p(s) (1+V(s)) ds dt \\ &\leq \bar{u}(r_0) + M \int_{r_0}^r e^{-t} t^{1-N} (1+V(t)) \int_0^t e^s s^{N-1} p(s) ds dt \\ &\leq \bar{u}(r_0) + M (1+V(r)) \int_{r_0}^r e^{-t} t^{1-N} \int_0^t e^s s^{N-1} p(s) ds dt, \end{split}$$

that is

$$\bar{u}(r) \le \bar{u}(r_0) + M(1 + V(r)) \int_{r_0}^r P(t) dt$$
, for all $r \ge r_0 \ge 0$. (16)

Since $\int_{1}^{\infty} P(r)dr < \infty$ and $\int_{1}^{\infty} Q(r)dr < \infty$, we can choose $r_0 \ge 1$ such that $\max\left\{\int_{r_0}^{\infty} P(r)dr, \int_{r_0}^{\infty} Q(r)dr\right\} < \frac{1}{2M}.$

From (14) and the fact that $\lim_{r \to \infty} \bar{u}(r) = \lim_{r \to \infty} \bar{v}(r) = \infty$, we can find $r_1 \ge r_0$ such that

$$U(r) = \max_{r_0 \le t \le r} \bar{u}(r), \quad V(r) = \max_{r_0 \le t \le r} \bar{v}(r), \quad \text{for all } r \ge r_1.$$
(18)

Thus (16) and (18) yield

$$U(r) \le \bar{u}(r_0) + M(1 + V(r)) \int_{r_0}^r P(t) dt$$
, for all $r \ge r_1$.

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(17)

Furthermore, by (17) we obtain

$$U(r) \le \bar{u}(r_0) + \frac{1 + V(r)}{2}$$
 for all $r \ge r_1$,

and so

$$U(r) \le C_1 + \frac{1}{2}V(r) \quad \text{for all} \ r \ge r_1,$$
(19)

where $C_1 = \frac{1}{2} + \bar{u}(r_0) > 0$. In a similar way we get

$$V(r) \le C_2 + \frac{1}{2}U(r)$$
 for all $r \ge r_1$, (20)

By addition, (19) and (20) lead to

$$U(r) + V(r) \le 2(C_1 + C_2)$$
 for all $r \ge r_1$. (21)

This means that U and V are bounded and so u and v are bounded which is a contradiction. It follows that (2) has no positive entire large solutions and the proof is now complete.

4. Proof of Theorem 3

We start by showing that (2) has positive radial solutions. On this purpose we fix a > 0 and b > 0 and we show that the system

$$\begin{cases} u'' + \frac{N-1}{r}u' + u' = p(r)f(v(r)), & r > 0, \\ v'' + \frac{N-1}{r}u' + v' = q(r)g(u(r)), & r > 0, \\ u', v' \ge 0 \quad \text{on } [0, \infty), \\ u(0) = a > 0, \ v(0) = b > 0, \end{cases}$$
(22)

has solutions. Then U(x) = u(|x|), V(x) = v(|x|) are positive solutions of (2). Integrating (22) we have

$$u(r) = a + \int_0^r e^{-t} t^{1-N} \int_0^t e^s s^{N-1} p(s) f(v(s)) ds dt \quad \forall r \ge 0,$$
(23)

$$v(r) = b + \int_0^r e^{-t} t^{1-N} \int_0^t e^s s^{N-1} q(s) g(u(s)) ds dt \quad \forall r \ge 0.$$
(24)

Define $v_0 \equiv b$ and let $(u_k)_{k \geq 1}, (v_k)_{k \geq 1}$ given by

$$u_k(r) = a + \int_0^r e^{-t} t^{1-N} \int_0^t e^s s^{N-1} p(s) f(v_{k-1}(s)) ds dt \quad \forall r \ge 0,$$
(25)

$$v_k(r) = b + \int_0^r e^{-t} t^{1-N} \int_0^t e^s s^{N-1} q(s) g(u_k(s)) ds dt \quad \forall r \ge 0.$$
(26)

Since $v_1(r) \ge b$, it follows that $u_2(r) \ge u_1(r)$ for all $r \ge 0$ which yields $v_2(r) \ge v_1(r)$ and so $u_3(r) \ge u_2(r)$ for all $r \ge 0$. Repeating such arguments we deduce that

$$u_k(r) \le u_{k+1}(r)$$
 and $v_k(r) \le v_{k+1}(r)$, for all $r > 0, k \ge 1$.

Let us now prove that the non-decreasing sequences $(u_k)_{k\geq 1}$ and $(v_k)_{k\geq 1}$ are bounded from above on bounded sets. We first observe that (25) and (26) yield

$$u_k(r) \le u_{k+1}(r) \le a + f(v_k(r)) \int_0^r P(t) dt, \quad \forall r \ge 0, \, k \ge 1$$
 (27)

and

$$v_k(r) \le b + g(u_k(r)) \int_0^r Q(t) dt, \quad \forall r \ge 0, \, k \ge 1$$
 (28)

Let R > 0 be arbitrary. From (27) and (28) we get

$$u_k(R) \le a + f\left(b + g(u_k(R))\int_0^R Q(t)dt\right)\int_0^R P(t)dt, \quad \forall k \ge 1.$$

This imply

$$1 \le \frac{a}{u_k(R)} + \frac{f\left(b + g(u_k(R))\int_0^R Q(t)dt\right)}{u_k(R)} \int_0^R P(t)dt, \quad \forall k \ge 1.$$
(29)

Taking into account the monotonicity of $(u_k(R))_{k\geq 1}$, there exists $L(R) := \lim_{k\to\infty} u_k(R)$.

We claim that L(R) is finite. Indeed, if not, we let $k \to \infty$ in (29) and the assumption (A_2) leads us to a contradiction. Thus L(R) is finite. Since u_k, v_k are increasing functions, it follows that the map $(0, \infty) \ni R \longmapsto L(R)$ is non-decreasing on $(0, \infty)$ and

$$\begin{split} u_k(r) &\leq u_k(R) \leq L(R), \quad \forall r \in [0, R], \, \forall k \leq 1, \\ v_k(r) &\leq b + g(L(R)) \int_0^R Q(t) dt, \quad \forall r \in [0, R], \, \forall k \leq 1. \end{split}$$

Furthermore, there exists $\lim_{R\to\infty} L(R) = \overline{L} \in (0,\infty]$ and the sequences $(u_k)_{k\geq 1}$ and $(v_k)_{k\geq 1}$ are bounded from above on bounded sets.

Let $u(r) := \lim_{k\to\infty} u_k(r)$, $v(r) := \lim_{k\to\infty} v_k(r)$ for all $r \ge 0$. By standard elliptic regularity theory we deduce that (u, v) is a positive solution of (22).

In order to conclude the proof, it is enough to show that (u, v) is a large solution of (22). Let us remark that (23), (24) imply

$$u(r) \ge a + f(b) \int_0^r P(t)dt, \quad \forall r \ge 0,$$
$$v(r) \ge b + g(a) \int_0^r Q(t)dt, \quad \forall r \ge 0.$$

Since f, g are positive functions and p, q satisfy (5) we can conclude that (u, v) is a large solution of (22) and so (U, V) is a positive entire large solution of (2). Hence any large solution of (22) provides a positive entire large solution (U, V) of (2) with U(0) = a and V(0) = b. Since $(a, b) \in (0, \infty) \times (0, \infty)$ was chosen arbitrarily, it follows that (2) has infinitely many positive entire large solutions. The proof of theorem is now complete.

Remark 4 The condition (5) is sufficient but not necessary for the existence of positive entire large solutions for (2). Indeed, let us consider $f(t) = \sqrt{t}$, g(t) = t, $p(r) = 4 \frac{r^3 + (N+2)r^2}{\sqrt{r^2+1}}$, $q(r) = 2 \frac{r+N}{r^4+1}$.

Using (6) we get $\int_{1}^{\infty} P(r)dr = +\infty$ and $\int_{1}^{\infty} Q(r)dr < +\infty$. However, the corresponding system to (2) is $\int \Delta u + |\nabla u| = 4 \frac{|x|^3 + (N+2)|x|^2}{\sqrt{v}} \cdot \sqrt{v} \quad \text{in } \mathbb{R}^N,$

$$\begin{aligned} \Delta u + |\nabla u| &= 4 \, \frac{|x|^3 + (N+2)|x|^2}{\sqrt{|x|^2 + 1}} \cdot \sqrt{v} & \text{ in } \mathbb{R}^N, \\ \Delta v + |\nabla v| &= 2 \, \frac{|x| + N}{|x|^4 + 1} \cdot u & \text{ in } \mathbb{R}^N, \end{aligned}$$

which has the positive entire large solution $(|x|^4 + 1, |x|^2 + 1)$.

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M. Ghergu and V. Rădulescu Department of Mathematics University of Craiova 200585 Craiova, Romania radulescu@inf.ucv.ro